# Effects of the Voltage and Pressure on the Carburizing of Martensitic Stainless Steel in Pulsed DC Glow Discharge

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In the present work, the influence of the pulse voltage and pressure on the treatment glow discharge characteristics and consequently on the surface properties obtained for low temperature plasma carburized AISI 420 martensitic stainless steel was investigated. Two distinct sets of samples were carburized at 450 °C, for 8 h, one aiming to study the applied pulse voltage effects, which was varied for 500, 600 and 700 V, at a fixed pressure of 400 Pa, and the other aiming to study the pressure effects, which was varied for 200, 400, and 800 Pa, at a fixed pulse voltage of 700 V. Treated samples were characterized by means of confocal laser scanning microscopy (CLSM), X-ray diffraction (XRD) analysis, microhardness and roughness measurements. The glow discharge (plasma) was characterized by optical emission spectroscopy (OES) and current measurements. Results show that the edge effect, surface roughness, hardness and outer layer growth kinetics are dependent on the studied plasma parameters. OES analyses showed that the pulse voltage parameter does not promote significant changes on the plasma chemistry, but confirmed that the molecular  $H_2$  gas dissociation rate tends to be significantly affected by the pressure parameter giving important support for the results obtained here.

**Keywords:** AISI 420 martensitic stainless steel, low-temperature pulsed dc plasma carburizing, pressure and pulse voltage parameters, optical emission spectroscopy, glow discharge (plasma) diagnostic.

## 1. Introduction

Martensitic stainless steels (MSSs) are widely used in different industrial fields comprising the manufacturing of surgical instruments<sup>1.3</sup>, forming tools<sup>3-5</sup>, hydraulic system elements<sup>3.5,6</sup>, structural and machine parts<sup>3,5,7</sup>, dairy and food apparatus<sup>1,4,7</sup>, petrochemical<sup>8</sup>, and mining industry<sup>9</sup> components. They present wide application field due to their relatively high mechanical strength and moderate corrosion resistance<sup>10</sup>. Nevertheless, application limitations are mainly found on conditions where high wear and corrosion resistances are required, in addition to adequate mechanical properties<sup>11</sup>.

Low temperature plasma carburizing (LTPC) treatment is a valuable surface engineering technique aiming to improve the MSSs corrosion performance<sup>12</sup> and their tribological response for specific systems<sup>13,14</sup>. LTPC means that the treatment temperature is high enough to promote the carbon interstitial diffusion, but low enough to avoid the chromium substitutional diffusion, resulting in *paraequilibrium* condition<sup>15</sup>. In the practical case, it occurs for treatment temperatures equal or lower than 450 °C and for times relatively short<sup>12-14,16</sup>. The obtained MSSs LTPC surfaces are usually constituted of cementite (Fe<sub>3</sub>C) and C-expanded martensite ( $\alpha'_{C}$ ) phases<sup>17</sup>, conferring significant strengthening as well as good corrosion resistance<sup>12</sup> for the treated surface. Differently, at high-temperatures that for MSSs carburizing means to be higher than 450 °C the formation of M<sub>23</sub>C<sub>6</sub>, M<sub>6</sub>C, M<sub>7</sub>C<sub>3</sub> phases provokes sensitization of the steel<sup>18</sup>.

Concerning the influence of the processing parameters studied here, the literature is rich for austenitic stainless steels (ASSs) plasma nitriding<sup>19-29</sup>, but very poor for MSSs, being mainly focused on treated surface characteristics of different materials, through different processes<sup>30-40</sup>. It is to be noted that almost nothing has been published aiming to study such effects on the MSSs LTPC. For ASSs, controversial results by comparing ref<sup>19,20</sup>. and ref<sup>21-24</sup>. have been reported for the pressure effects on the layer hardness and thickness, but all these authors agree very well that the pressure increment interestingly promotes an increase in the treated material corrosion resistance<sup>20,22-24</sup>, possibly due to the combined effect of the Cr kept in the steel matrix solid solution and the protective layer formation<sup>23,25</sup>. Regarding the plasma, the pressure increase promotes an increase in electron and atomic nitrogen densities26. Differently, only one work was

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found on the pulse voltage effect, which showed slight layer hardness increase for an ASS<sup>28</sup>.

Considering the carburizing treatment, authors of this work have previously studied the influence of the carburizing gas mixture and flow rate<sup>41</sup>, and the carburizing temperature and time<sup>17</sup> on the layer formation in the MSS. So, in this work, the influence of the pulse voltage and the (gas) pressure on the MSS carburizing process carried out in a pulsed direct current (*dc*) abnormal glow discharge is investigated. The main goal of this work is to contribute to the understanding of how these two carburizing glow discharge parameters and the respective plasma constitution can be directly related to the carburized layer formation process in the present treatment.

#### 2. Bases for the Carburizing Treatment in *DC* Glow Discharge

In the LTPC process, time is an independent parameter, differently from the temperature, which is strongly dependent on the glow discharge and/or electrical parameters. Regarding the pulsed dc plasma thermochemical treatments, pulsed wave (voltage and duty cycle) parameters and pressure strongly affect the number density and energy of plasma species<sup>42-44</sup>. Such parameters are the main responsible to control the treatment temperature of substrates or parts acting as cathode in the abnormal glow discharge regime, which depends on the plasma species bombardment<sup>45</sup>. They also present significant technological importance. When separately working, pressure influences the species mean free path, while the applied pulse voltage influences the average electrons energy, plasma density and energy of ions bombarding the cathode surface<sup>42,46,47</sup>. Differently, considering these two parameters working together, for a dc plasma process, the Davis and Vanderslice (1963) rule is supposed to operate. This rule predicts that the product between pressure (p) and cathode sheath thickness (d) tends to be constant for a specific voltage kept unaltered<sup>48</sup>. It means that for a constant voltage, the effect of pressure on the ions energy distribution is small, since rising or reducing the pressure implies in reducing or rising the cathode sheath thickness, respectively, so that a relatively constant number of collisions in the cathode sheath, which directly infer changes on the species energy, can be expected. Differently, for a constant pressure, the increase of the voltage promotes an increase on the cathode fall and a respective reduction on the cathode sheath thickness, thus the ions average energy also increases48. Such aspects associated with the inelastic collision processes<sup>49</sup> leading reactive species to be formed summarize how the reactive species density directly influence the physical-chemical interactions in the plasma-surface interface, thus on the thermochemical treatment result. In this case, it is worth mentioning different phenomena occurring such as the edge effect<sup>19,20</sup>, sputtering<sup>42-44,50,51</sup>, active species generation<sup>52,53</sup>, and the cathode (samples) heating, which in the practice is a function of the pulse switched-on time or duty cycle  $(t_{op})^{51}$ .

#### **3.** Experimental Procedure

#### 3.1. Materials

Cylindrical samples of AISI 420 steel (for a composition of 0.17% C, 0.70% Mn, 0.50% Si, 12.22% Cr, 0.16% Ni,

0.23% P, 0.03% S, 0.03% N, 0.01% Cu, 0.02% Co, and balance in Fe, in wt.%) were cut to 10 mm in height from a 9.5 mm diameter commercial rod supplied in annealed state. In order to obtain a fully martensitic structure<sup>54-56</sup>, samples were austenitized at 1050 °C for 0.5 h and oil quenched, reaching in this as-quenched state an average hardness of  $510 \pm 10 \text{ HV}_{0.3}$ . Before LTPC treatment, samples were ground using SiC sandpaper up to 1200 grade and mirror polished using 1  $\mu$ m Al<sub>2</sub>O<sub>3</sub> suspension, which resulted for average *Ra* and *Rz* roughness values of 0.070 ( $\pm$  0.008) and 0.320 ( $\pm$  0.027)  $\mu$ m, respectively. Afterwards, samples were cleaned in ultrasonic bath using isopropyl alcohol, dried under heated airflow and introduced into the glow discharge chamber.

## 3.2. The LTPC processing

Scheme and complete description of the plasma treatment apparatus utilized in this study can be found in ref57. After evacuating the system to a pressure of 4 Pa (30 mTorr), samples were subjected to (plasma) sputter-cleaning at 300 °C, for 0.5 h, in an 80% H<sub>2</sub> + 20% Ar gas mixture, for flow rate of  $1.67 \times 10^{-6} \text{ Nm}^3 \text{.s}^{-1}$  (100 sccm), under 400 Pa (3 Torr) pressure. As-quenched samples were carburized at 450 °C, for 8 h, in 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture, keeping the same pressure and flow rate used in the previous (plasma) sputter-cleaning step. In the present study, as-quenched samples were simultaneously tempered during the plasma carburizing treatment, as expected from the thermal effect; being of interest details about the steel tempering can be found in ref<sup>58</sup>. Once concluded the carburizing step samples were cooled down to room temperature under  $H_{2}$  + Ar gas mixture flow.

Two treatment series using a 4.16 kHz square-wave pulsed dc power supply and a pulse period of 240 µs were carried out:

- i. one, varying the pulse voltage parameter for 500, 600, and 700 V at a fixed pressure of 400 Pa. For these treatments the  $t_{on}$  (the pulse switched-on time or duty cycle) was 136 (±4), 70 (±1) and 47 (±1) µs, and the electrical current was 166 (±3), 144 (±3) and 133 (±3) mA, respectively; and
- ii. the other, varying the pressure for 200, 400, and 800 Pa (1.5, 3 and 6 Torr, respectively) at a fixed pulse voltage of 700 V (actual 688 ± 12 V). For these treatments the  $t_{on}$  was 128 (±2), 50 (±1) and 50 (±1) µs, and the electrical current was 126 (±7), 137 (±2) and 169 (±2) mA, respectively.

The observed  $t_{on}$  values needed to attain the treatment temperature and the current dependence on this independent parameter, namely the  $t_{on}$ , is explained by the LTPC working in the abnormal glow discharge regime, as observed in ref<sup>50,51,59,60</sup>.

#### 3.3. The glow discharge characterization

In order to qualitatively evaluate the emitting plasma species as a function of the applied pulse voltage and pressure, and aiming at relating the different obtained plasmas with the respective treated surface layers characteristics, experiments were separately conducted for optical emission spectroscopy (OES) measurements. For the purpose to avoid possible variations in experimental setup, which would forbid line intensity comparisons, OES measurement experiments were not carried out during the actual treatments. For this case, measurements were performed using an HR4000 Ocean Optics Spectrometer equipped with a 3648-element linear CCD array and a 300 lines/mm grating, set to operate in the range of 200-1100 nm, with a wavelength resolution of 1.3 nm FWHM (full width at half maximum). In the related experiments, for the first series, nominal applied pulse voltage was changed starting from 500 V, and adding 100 V each 60 min. For the second series, the pressure was initiated at 200 Pa, and it was changed to 400, and after to 800 Pa, keeping a 60 min permanence time for each adjusted pressure. In both cases, optical spectra were collected from the beginning of the parameter setting (pressure or pulse voltage) until the end of the 60 min, even if the steady state in OES was attained in the first 5 min. Each spectrum was collected as an average of 30 spectra of 1 second of exposure time, in steady state operation. During such experiments, the  $t_{ar}$  was adjusted as a function of the applied pulse voltage and pressure, in order to maintain the temperature at 450°C. It is to be noted that the experiment was designed this way aiming to ensure the same measurement solid angle (optical fiber was fixed in the same position all long the measurement series), so the emission spectra for each applied pulse voltage and pressure condition of interest was acquired without any changes in the OES experimental setup. As the position between sample and optical fiber was kept unaltered, the relative intensity of the obtained emission lines, of the normalized spectra, can be confronted.

#### 3.4. The treated samples characterization

Carburized samples were cross-sectioned and prepared for microstructural characterization according to conventional metallographic preparation procedures. Villela's reagent (95 ml of ethyl alcohol, 5 ml of hydrochloric acid, and 1 g of picric acid) was used to reveal the microstructure. Microstructural analysis was performed employing a Confocal Laser Scanning Microscope (Olympus LEXT OLS 3000); this same equipment was used to determine the average Ra and Rz roughness of treated surfaces, but in this case using a 200x magnification, applying low-pass filtering and slices number suggested by the software dedicated to the microscope control. The determination of carburized layer phases by Xray diffraction (XRD) was carried out as previously indicated in ref<sup>16</sup>. Vickers (HV) and Knoop (HK) microhardness measurements were also performed applying a load of 10 gf during 15 s. For the former, indentations were made in the central part of carburized (top) and non-carburized (bottom) flat surfaces, at a distance of 100 µm from each other. For the latter, the respective hardness profiles were obtained, being the measurement distance along the depth fixed in 10 µm. All indicated hardness values correspond to the average and standard deviation of five measurements performed for each indicated point, being that in graphics, each point presents error bars representing the dispersion of the measured values around the mean value. Finally, the carburized layer depth was determined by means of hardness profiles, using a similar procedure as proposed in NHT standard (DIN 50190-3:1979-03 used in nitriding), which consists in finding the depth where the hardness is 50 HV higher than that of the substrate bulk, but for the present case it was used 50 HK as reference.

#### 4. Results and Discussion

# 4.1. Microstructural aspects and treated surface characteristics

All carburized samples showed a predominantly-whiteaspect layer at the treated surfaces. Two typical microstructures of the applied pulse voltage and pressure treatment series are presented in Figure 1a, b, one obtained at 500 V / 400 Pa and the other at 800 Pa / 700 V, just the lower voltage and the higher pressure studied here, respectively. This result clearly indicates that the carburized layer it is more chemically resistant to the etching with Villela's reagent than the sample substrate bulk. The white-aspect layer occurrence suggests carbon diffusing and enriching in interstitial solid solution



Figure 1. Typical cross-section microstructures of the applied pulse voltage and pressure treatment series obtained at: a) 500 V / 400 Pa; and b) 800 Pa / 700 V, respectively. Treatments carried out at 450 °C, for 8 h, using  $99.5\% (80\% \text{ H}_2 + 20\% \text{ Ar}) + 0.5\% \text{ CH}_4$  gas mixture at  $1.67 \times 10^{-6} \text{ Nm}^3 \text{ s}^{-1}$  flow rate.

the surface martensitic structure, promoting the formation of the termed 'outer layer', according to the terminology previously used in ref<sup>41</sup>. The occurrence of diffusion layer, mainly formed by the carbon interstitial solid solution into the substrate bulk cannot be visualized in the micrographs, being only identified through the hardness profiles, as presented ahead.

Figure 2a, b left y-axis shows Vickers hardness values measured on the carburized (top) and non-carburized (bottom) flat surfaces, obtained for the applied pulse voltage and pressure treatment series, respectively. The top surface hardness grows slightly with increasing of the applied pulse voltage, exhibiting values of  $912 \pm 24$ ,  $975 \pm 26$  and  $1022 \pm 23$  HV<sub>0.01</sub> for samples carburized at 500, 600 and 700 V, respectively. The hardness values measured on the bottom surface, in this order, were  $392 \pm 28$ ,  $387 \pm 33$ , and  $402 \pm 28$  HV<sub>0.01</sub>. Comparatively, results indicate a hardness increase on the order of 250% for the carburized surfaces, being slightly lower at 500 V, and a hardness decrease on the order of 20% for all non-carburized surfaces due to the steel tempering effect, a result valid for both the studied series. Differently, the carburized surface hardness shows to be sensitive to the pressure parameter. In this case, by varying the pressure from 200 to 400 Pa, a considerable hardness increase from  $813 \pm 36$  to  $1022 \pm 23$  HV<sub>0.01</sub> was verified, the same being not valid by changing it from 400 to 800 Pa, which hardness slightly decreases from  $1022 \pm 23$  and  $966 \pm 25$  HV<sub>0.01</sub>, respectively.

Figure 2a, b right y-axis shows the obtained outer layer thickness and average treated surface Ra and Rz roughness variation. The layer thickness data demonstrate similar behavior to those obtained for the measured hardness, confirming that thicker layers tend to be harder than the thinner ones. In this case, a slight increase ranging from 1.9 to 2.2 µm was observed by increasing the pulse voltage from 500 to 700 V. The higher pulse voltage the higher is the average plasma species<sup>†</sup> bombardment energy, favoring the MSS oxide layer destabilization and sputtering of the surface atoms. Note that the mean electrical current needed to achieve the treatment temperature decreased from 166 mA (for  $t_{ov} = 136 \ \mu s$ ) to 133 mA (for  $t_{\rm er} = 47 \,\mu s$ ) when the pulse voltage was increased from 500 to 700 V, respectively. As previously indicated, the measured electrical parameters variation is due to the fact that the treatment has occurred in the abnormal glow discharge regime, in which the current approximately vary linearly with the applied voltage, also being a function of the duty cycle  $(t_{av})$  for the case of a pulsed dc glow discharge. In other words, this explains why for the 700 V voltage condition was needed a  $t_{on}$  value about three times smaller than that observed at 500 V voltage. Finally, since the cathode temperature is dependent on the ion current density and the thermal losses to the surrounding walls (from the grounded anode, which is constituted by thermal shields and discharge chamber walls), this explains the observed roughly similar current values independently of the utilized pulse voltage.

In addition, the pulse voltage increment also intensifies the electron impact ionization collisions, thus the molecules species dissociation rate. Such effects tend to increase the

carbon precursor active species concentration and atomic carbon available on the treating surface, which is able to diffuse into samples. All these combined effects allied to the fact that the surface carbon enrichment in this LTPC process is diffusion-controlled / -limited16 help to explain the slight increase on both the carburized layers thickness (~0.3 µm) and hardness ( $< 110 \text{ HV}_{0.01}$ ). Similarly, the above-presented relation between the layer thickness and its hardness was also observed for the pressure treatment series, which showed a 1.7 µm thinner layer at 200 Pa, and approximately equal thicker layers for samples treated at 400 and 800 Pa, both close to 2.2 µm thickness. This result also indirectly indicates a possible increase on the active species and plasma densities. Additionally, pressure also influences the plasma species collision frequency and the mean free path of gas particles<sup>42,43,47,49</sup>. This means that one can expect a greater frequency of collisional processes, influencing the formation of CH<sub>i</sub> (i = 0, 1, 2, 3) reactive species, thus enhancing the carbon-based reactive species production from CH<sub>4</sub> molecules<sup>61</sup>. On the other hand, the shortening of the mean free path at higher pressures will cool down electrons, and, for too high pressure, the dissociative collisions effectiveness can decrease. Based on this reasoning, the relatively small outer layer thickness and hardness for samples treated at 200 Pa can be justified by the lower carbon atoms density supplied by the plasma to the treating surface. Furthermore, the similarity of the findings for the conditions of 400 and 800 Pa is possibly justified by the fact that for the higher treatment pressures, the carbon supplying supposedly was not altered significantly. Such results mean a condition for which an active species yielding steady state was reached, probably associated to the eventual thermalization or, differently, the carbon supplying was greater than the capacity of the material to absorb it, strongly indicating that the treatment became a diffusion-limited process, as previously already indicated here. Finally, another and very probable assumption to explain the difference obtained at 200 Pa, regarding the similar results at 400 and 800 Pa, is the well-known role of the atomic hydrogen in promoting decarburization, as discussed from OES results, presented ahead.

Figure 2a, b right y-axis also shows the Rz and Ra roughness of the treated surfaces. Results indicate a slight roughness increase tendency as the pulse voltage is increased. Increasing roughness values of 0.75, 0.98, and 1.0 µm for Rz, and 0.12, 0.15 and 0.18 µm for Ra were observed for samples carburized at 500, 600 and 700 V, showing significant changes when compared to the starting 0.320 µm Rz and 0.070  $\mu m Ra$  values, respectively. This behavior can be explained by the fact that the energy for which the ions and fast neutrals collide with the substrate surface is directly related to the power transferred to the plasma by the applied voltage42,46,47,49, leading the sputtering effect to be also increased. Results also demonstrated that the Rz and Ra roughness significantly decreases as the pressure is increased from 200 to 400 and 800 Pa, being obtained values of 1.2, 1.0 and 0.5 µm for Rz and 0.25, 0.23 and 0.18 µm for Ra, respectively. It is known that the sputtering may promote changes on the samples surface morphology, being strongly dependent on the pressure<sup>42</sup>. The smaller the pressure the higher is its effect at the treated surface, remembering that the

<sup>†</sup> ions and fast neutrals produced from symmetrical and asymmetrical charge change collisions in the cathode sheath<sup>38-49</sup>.

backscattering phenomenon is enhanced for higher pressures<sup>49</sup>. This means that a larger sputtered atoms density diffuses back to the parent surface, decreasing the sputtering effect on the observed roughness. To the case of the lower pressure (200 Pa), the larger mean free path may lead sputtered atoms after thermalization to diffuse away from the parent surface, due to the lower backscattering effect, thus resulting in the increment of the surface roughness, as observed in ref<sup>42,49</sup>. In brief, the resulting higher sputtered atoms redeposition rate for the highest studied pressure (800 Pa) associated to the presumable redeposition of these atoms on valleys of the parent surface would be explain the observed Rz and Ra roughness reduction. Finally, the electrical parameters variation from 126 mA (for  $t_{\rm er} = 128 \,\mu s$ ) at 200 Pa to 169 mA (for  $t_{av} = 50 \ \mu s$ ) at 800 Pa, needed to keep unaltered the treatment temperature, would agree with the higher density of lower-energy ions bombarding the cathode surface at the higher pressure. This increased lower-energy ion density, expected at the higher pressure treatment (800 Pa) and directly resulting in a higher current seems to compensate in terms of the cathode heating effect the relatively reduced higher-energy ion density, leading to a smaller current at

400 Pa treatment pressure, so explaining the same  $t_{on}$  value of 50 µs observed for both the 400 and 800 Pa treatments, as confirmed by the measured electrical current for each case, which was 137 (±2) and 169 (±2) mA, respectively.

Figure 3a, b shows the Knoop hardness profiles measured on the carburized samples, obtained for the applied pulse voltage and pressure treatment series, respectively. The 2.5 µm distance from the surface, observed for the first measurement performed in the samples cross-section, is greater than the obtained outer layers thickness. To clarify, the outer layer region (constituted of  $\alpha'_{C}$  + Fe<sub>3</sub>C phase) is also indicated in the Figure 3a, b detail. In addition, all hardness profiles confirm the existence of typical diffusion layers showing a smooth hardness decrease from the treated surface to the tempered substrate bulk. Carburized layer depth on the ~40-45 µm range was obtained for all studied conditions, with the exception of that obtained at 200 Pa, which presented ~25 µm depth and, comparatively, the smaller hardness value at the 2.5  $\mu$ m depth. Hardness values of 845  $\pm$  16, 877  $\pm$  11 and  $882 \pm 12$  HK<sub>0.01</sub> for samples carburized at 500, 600, and 700 V, and 643  $\pm$  15, 860  $\pm$  12 and 812  $\pm$  10  $\rm HK_{0.01}$  for



Figure 2. Vickers hardness measured on the carburized (top) and non-carburized (bottom) flat surfaces, layer thickness, and surface *Ra* and *Rz* roughness obtained for: a) applied pulse voltage; and b) pressure treatment series. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at  $1.67 \times 10^{-6}$  Nm<sup>3</sup> s<sup>-1</sup> flow rate.



Figure 3. Knoop hardness profiles measured on the carburized samples obtained for: a) applied pulse voltage; and b) pressure treatment series. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at  $1.67 \times 10^{-6}$  Nm<sup>3</sup> s<sup>-1</sup> flow rate.

samples treated at 200, 400 and 800 Pa, were measured to the 2.5  $\mu$ m depth, respectively.

Figure 4a, b shows X-ray diffraction patterns of untreated and carburized surfaces, for the applied pulse voltage and pressure treatment series, respectively. All results confirm the obtainment of the carbon-expanded martensite ( $\alpha'_{c}$ ) phase, since the indicated peaks cannot be matched with any other possible phases, including iron and chromium carbides. In addition, the martensite  $(\alpha')$  peak initially presented in the untreated sample were broadened and shifted to lower diffraction angles in the carburized samples, indicating lattice parameter expansion by the carbon and residual stress formation, as previously obtained in ref12-14,28,62. for LTPC, and in ref63. and 64, for low-temperature solid carburizing and low-pressure carburizing (LPC) techniques, respectively. The  $\alpha'_{C}$  phase peak for samples treated at 600 and 700 V were displaced to smaller angles when compared to that at 500 V. This suggests greater carbon supersaturation for those conditions, being in agreement with the related Figure 2 and 3 results. Moreover, all XRD patterns obtained for samples of both the studied treatment series also show peaks at 39.8° and 45.9° that correspond to the cementite (Fe<sub>2</sub>C) phase (according to the JCPDS card n. 34-0001). Cementite or Fe<sub>2</sub>C phase (in fact, M<sub>2</sub>C-type carbide, for which Cr would partially substitute Fe atoms) peaks show greater intensities for samples treated at 600 and 700 V, also suggesting the formation of higher iron-based carbide phase volumetric fractions. Regarding the pressure treatment series, changes on the original martensite ( $\alpha$ ') peak for the sample treated at 200 Pa are less pronounced than those observed for samples treated at 400 and 800 Pa. This result is related to the possible lower carbon supply that is in agreement with the lower reactive species density for this sample, as previously discussed, disregarding the XRD probing depth, since for this condition (200 Pa) the outer layer is thinner, so that the XRD pattern would have a higher contribution coming from the diffusion layer and lower information coming from the outer layer, when compared to the other studied pressures. Despite the formation of Fe<sub>2</sub>C phase in

stainless steels is not common, since chromium carbides are much more stable, its obtainment in the studied samples is due to the LTPC treatment paraequilibrium conditions. As discussed in details in ref<sup>65</sup>, for treatments at relatively low temperatures ( $\leq 450$  °C) the Fe and Cr (substitutional elements) diffusivity in the stainless steel matrix is negligible whereas that of the C (interstitial element) is significant. Since Fe atoms are much more available in a stainless steel alloy than Cr ones, as expected for a ferrous alloy, the probability for iron carbide phase formation is much higher at paraequilibrium. In other words, chromium carbide phases preferentially tend to be formed at temperatures higher than 450 °C and/or treatment times longer than 12 h, when the Cr diffusivity becomes significant, as shown in ref<sup>16</sup>.

To conclude this point, different aspects of the treated surface morphology are shown in Figure 5a, b. Despite it is very well known on the plasma nitriding field<sup>66,67</sup>, no work on edge effect from the plasma carburizing treatment was found in literature. Figure 5a shows the edge effect aspect and hardness variation along the treated surface for samples of the pressure treatment series. This effect presents significant technological interest for *dc* plasma treatments, usually leading to heterogeneous properties and characteristics varying from the sample or part boundary towards its center at the treated surface, such as those verified here.

So, hardness measurements spaced for  $250 \,\mu\text{m}$  one each other and radially made from the surface into the sample center showed hardness decrease on the order of 50 HV for the studied samples (see Figure 5a), as also shown in ref<sup>27,28</sup>. This result is associated to the higher thermal effect located together the sample boundary associated to a more intense sputtering, suggesting differentiated and located pulverization of carbon atoms condensed at the treating surface. Besides the reduced carbon enrichment in such boundary regions, the higher plasma density due to the sum of two glow regions, one formed on the sample flat surface and the other formed around the sample cylindrical surface, tends to intensify a located tempering effect, thus also contributing to the reduction of its hardness. Figure 5a also shows the



**Figure 4.** X-ray diffraction patterns of untreated and carburized surfaces obtained for: a) applied pulse voltage; and b) pressure treatment series. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at  $1.67 \times 10^{-6}$  Nm<sup>3</sup> s<sup>-1</sup> flow rate.

flat surface of treated specimens clearly indicating circular marks varying on the order of 1.0, 0.5 and 0.25 mm were obtained at 200, 400 and 800 Pa, respectively, confirming that more heterogeneous-aspect surfaces are obtained for higher pressures, which tend to reduce this studied effect. So, the edge effect reduction is corroborated by more restricted glow regions that can be achieved using higher treatment applied pulse voltages and/or pressures, causing electric field distortions<sup>68</sup> only very close to the vicinity of sample boundaries, reducing the cathode sheath thickness,



**Figure 5.** (a) Edge effect aspect and hardness variation along samples of the pressure treatment series; and (b) Typical surface topography images of the sample treated at 800 Pa / 700 V. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at 1.67 × 10<sup>-6</sup> Nm<sup>3</sup> s<sup>-1</sup> flow rate.

and leading the plasma to more homogeneously cover each respective surface.

Finally, Figure 5b shows a typical surface morphology image for LTPC treated samples, in this case for the sample carburized at 800 Pa / 700 V. This morphology, revealed by sputtering, shows typical coarsened (plate-like) martensite<sup>65</sup>, agreeing very well with the significant increment observed for Rz and Ra roughness measurements of the treated samples.

# 4.2. Glow discharge characteristics used in the LTPC

Glow discharge spectra of treatments carried out at 500, 600 and 700 V are shown in Figure 6a. It is important



**Figure 6.** (a)  $t_{on}$ -normalized intensity OES spectra of the glow discharge; (b) instantaneous H<sub>a</sub> and Ar e.s. emission intensity (represented by the Intensity /  $t_{on}$  axis spectrum evolution) and H<sub>a</sub>/Ar e.s. intensity ratio; and (c) global intensity OES spectra (not divided by  $t_{on}$ ), for treatments carried out at 500, 600 and 700 V. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at 1.67 × 10<sup>-6</sup> Nm<sup>3</sup> s<sup>-1</sup> flow rate, and pressure of 400 Pa.

to observe that in order to maintain the same treatment temperature at a fixed pressure of 400 Pa, using different bias voltage, it was mandatory to adjust the  $t_{ov}$  (duty cycle) of the pulsed dc power supply. For this reason, this parameter needed to be took into account when the intensity of the emission lines was analyzed (represented by the Intensity /  $t_{an}$ ordinate axis in Figure 6a, b). This same observation is also valid for the pressure variation study, presented ahead (see Figure 7a, b). For the considered 300-900 nm wavelength range, as shown, the most intense detected emission lines are related to atomic Hydrogen  $(H_{\alpha}, H_{\beta})$  and Argon (Ar) emissive species (e.s.). Hydrogen lines, namely H<sub>a</sub> (656 nm) and  $H_{\alpha}$  (486 nm), and several lines of Ar in the range of 730-850 nm are the most intense emissions detected in the studied conditions. Despite the presence of C<sub>u</sub>H<sub>u</sub> species in the plasma, the relative intensity of emission lines from these species was not intense enough to be detected, which is related to the relatively low CH4 number density in the treatment gas mixture. As expected, as the applied voltage rises more intense instantaneous emissions intensity (represented by the Intensity /  $t_{av}$  axis) was verified. This is a result of an increase in the yield of excited (metastable) H and Ar species due to the increase of the electron temperature and number density with the increasing pulse voltage. Because of that, a greater instantaneous number density of active species is found, during the switched-on period, when the pulse voltage is augmented. In accordance with Figure 6b results, when the emission intensity of atomic Hydrogen (H) and Argon (Ar) emissive species (e.s.) are taken into account, a roughly stable H<sub>a</sub>/Ar e.s. intensity ratio is observed on the 500-700 V range. So, the applied pulse voltage seems to play no important role, despite the very slight observed growing tendency on the relative creation rate of H and Ar (e.s.), by increasing the voltage. This result is confirmed in Figure 6c by the approximately equal spectra in the three studied conditions when considering the global intensity (not divided by  $t_{\rm ex}$ ), indicating that the pulse voltage has similar influence on the creation mechanism of H and Ar (e.s.). Considering that such species are mainly created by electron impact, probably the same behavior is also valid for C<sub>v</sub>H<sub>v</sub> species. It suggests that in the studied pulse voltage range, the variation of this parameter more significantly modifies the electron number density than its energy distribution.

Glow discharge spectra of treatments carried out at 200, 400 and 800 Pa are shown in Figure 7a. Atomic Hydrogen (H) and Argon (Ar) emissive species (e.s.) were again the most intense emissions detected. However, differently of the results obtained from the pulse voltage variation, quite distinct spectra by increasing the treatment pressure from 200 to 800 Pa were obtained. As can be noted in Figure 7a, b, by varying the treatment pressure from 200 to 400 Pa, the spectra intensity is increased, but an intensity drop is observed by increasing the pressure from 400 to 800 Pa. Such drop should be related to the smaller mean free path at higher pressures, as also evidenced in results of the ref19, which would result in less energetic collisions, thus in a fall of the respective intensities. In addition, by analyzing the region from 640 to 840 nm (see Figure 7a), a more significant drop for  $H_{\alpha}$  than that of the Ar (e.s.) at 750 nm occurs when the pressure is increased (Figure 7c). This change in the relative intensity



**Figure 7.** (a)  $t_{on}$ -normalized intensity OES spectra of the glow discharge; (b) instantaneous H<sub>a</sub> and Ar e.s. emission intensity (represented by the Intensity /  $t_{on}$  axis spectrum evolution) and H<sub>a</sub>/Ar e.s. intensity ratio; and (c) global intensity OES spectra (not divided by  $t_{on}$ ), for treatments carried out at 200, 400 and 800 Pa. Treatments carried out at 450 °C, for 8 h, using 99.5% (80% H<sub>2</sub> + 20% Ar) + 0.5% CH<sub>4</sub> gas mixture at 1.67 × 10<sup>-6</sup> Nm<sup>3</sup> s<sup>-1</sup> flow rate, and voltage of 700 V.

ratio (Figure 7b) means that there is a considerable change in the plasma species energy distribution on the studied pressure range. Kobashi<sup>69</sup> stated that the H $\alpha$ /Ar (e.s.) intensity ratio provides an approximated measure for atomic hydrogen (H) number density in plasma. So, this result can bring an additional interpretation for the treated surface hardness and hardness profile variation with the treatment pressure. Atomic hydrogen (H) is known to promote decarburizing, so when carbon active species condense on and react with the surface, there is a competition between its diffusion towards the substrate bulk and the surface decarburizing. Thus, a high H number density in the plasma atmosphere can lead to a relatively lower actual carburizing potential along the treatment. Finally, considering obtained data of XRD, and the hardness and treated layer thickness characterization, it seems clear that the H number density plays an important role on the low-temperature plasma carburizing. For the studied conditions, higher pressures (400 and 800 Pa) were more effective to produce  $\alpha'_{\rm C}$  hard layers than the lower pressure (200 Pa), which agrees very well with its higher H number density in the plasma atmosphere (see Figure 7c). It is assumed that the higher H number density at 200 Pa is related to a relatively lower actual carburizing potential, as suggested from Figure 2, 3, and 4a, b results for the present case, as well as a higher dissociation rate of the molecular H<sub>2</sub> species.

### 5. Conclusions

A study strengthened by an optical emission spectroscopy approach for low-temperature plasma carburizing was performed aiming to evaluate the influence of the treatment pressure and applied pulse voltage on the glow-discharge characteristics, and on surface properties of the treated AISI 420 martensitic stainless steel. The main conclusions of this work are:

- The variation of the applied pulse voltage, in the studied range, does not significantly influence the surface characteristics of treated samples. The average values of outer layer hardness and thickness as well as the treated surface roughness grow modestly with the pulse voltage varying from 500 to 700 V;
- The variation of the treatment pressure leads to significant changes in the surface characteristics of the treated samples. Pressures of 400 and 800 Pa are responsible by the most promising results for the outer layer thickness and hardness. Likewise, surface roughness as well as edge effect decrease with the increment of the pressure; and
- Optical emission spectroscopy (OES) showed to be an important complementary technique to give support for the findings obtained in this work. Optical emissions spectra are in agreement with the metallurgical results that indicate more effective carburizing treatments occurring for pressures of 400 and 800 Pa. Results evidenced for the lower pressure (200 Pa) suggest the significant decarburizing role played by the atomic hydrogen. Comparatively, the dissociation rate of the molecular H<sub>2</sub> species tends to be significantly affected by the pressure parameter.

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